

Weak boson scattering from weak boson PDFs

Snowmass on the Zoomissippi: EF04 Group Meeting

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¹on behalf of many co-signers

the big *Lol* question

Question: How does the electroweak theory behave at very high energies?

Letter of Interest: EW effects in very high-energy phenomena

C. ARINA, G. CUOMO, T. HAN, Y. MA, F. MALTONI, A. MANOHAR, S. PRESTEL, R. RUIZ,
L. VECCHI, R. VERHEYEN, B. WEBBER, W. WAALEWIJN, A. WULZER, K. XIE
to be submitted to the Theory Frontier (TF07) and Energy Frontier (EF04)

1 Introduction

Phenomena that take place at multi-TeV scales — high-energy elementary particle scattering or the annihilation/decay of ultra heavy states such as dark matter particles — can give rise to relativistic, final states that are naturally accompanied by additional radiation, that in turn leads to particle showers and final states with large particle multiplicities. In the Standard Model, the effects of QCD and QED radiation are well understood and treated at various level of sophistication. These range from fixed-order computations at an increasing accuracy to resummed computations via parton showering algorithms and semi-analytic approaches. Even matching/merging between the two while keeping their respective accuracies is available.

In such multi-TeV scales processes typical momentum transfers Q are much larger than the electroweak (EW) scale $m \sim m_Z$, and initial- and final-state EW radiation becomes important. In particular, EW boson emission gives rise to transition rates that grow with logarithms of the type $\log Q/m$. For sufficiently large Q , these logarithms must be resummed in order to recover physically meaningful results. Despite recent progress, a fully exclusive approach that can take care of fixed-order EW corrections, resum large EW logarithms in both initial and final states, systematically account for power corrections, and is implemented in ready-to-use Monte Carlo

Snowmass 21 LoI: SNOWMASS21-TF7_TF0-EF4_EF0-026

Many *fascinating* ways to explore this, e.g., EW parton showers

Summary of output since hiatus

Lots of neat work from the past year! (lots of older work uncited)

- EW boson scattering / fusion at future colliders (**review**)

⌚ Buarque, et al [2106.01393]

- EW boson PDFs

Han, Ma, Xie [2007.14300; 2103.09844]; ⌚ Ruiz, et al [2111.02442]

- EW parton showers

Veryheyen, et al [2002.09248; 2108.10786]

- Spin correlation and interference in parton showers

Prestel, et al [2109.09706; 2110.05964], Veryheyen, et al [2111.01161]

- Muon collider physics

(many!)

How the EW theory behaves at very high energies is being explored!

- Exploring polarized matrix elements with MadGraph5

w/ D. Buarque Franzosi, O. Mattelaer, S. Shil [1912.01725]

- Exploring EW VBF with multi-TeV muon colliders

w/ A. Costantini, F. De Lillo, F. Maltoni, L. Mantani, O. Mattelaer, X. Zhao [2005.10289]

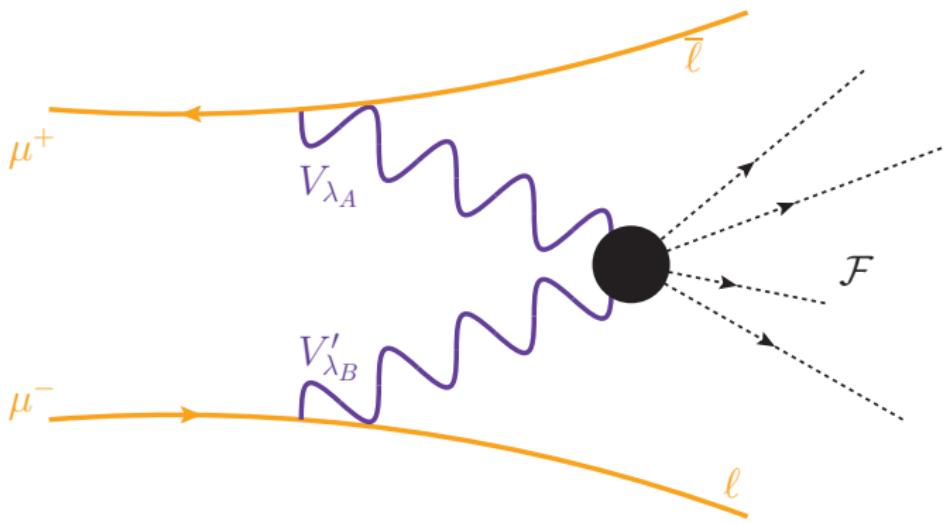
- ▶ Multi-TeV $\ell^+\ell^-$ colliders are “high-luminosity EW boson colliders”

- NEW! Exploring the validity of W_λ/Z_λ PDFs

w/ A. Costantini, F. Maltoni, O. Mattelaer [2111.02442]

- ▶ Support for W_λ/Z_λ PDFs from μ^\pm, e^\pm beams in MadGraph5

Part III: Effective Vector Boson Approximation in High-Energy Muon Collisions



What did we implement into MadGraph5_aMC@NLO?

$$\sigma(\mu^+ \mu^- \rightarrow \mathcal{F} + X) = f \otimes f \otimes \hat{\sigma} + \text{power-law and log corrections} \quad (1)$$

What did we implement into MadGraph5_aMC@NLO?

$$\sigma(\mu^+ \mu^- \rightarrow \mathcal{F} + X) = f \otimes f \otimes \hat{\sigma} + \text{power-law and log corrections} \quad (1)$$

$$\begin{aligned} &= \sum_{V_{\lambda_A}, V'_{\lambda_B}} \int_{\tau_0}^1 d\xi_1 \int_{\tau_0/\xi_1}^1 d\xi_2 \int dPS_n \\ &\times \left[f_{V_{\lambda_A}/\mu^+}(\xi_1, \mu_f) f_{V'_{\lambda_B/\mu^-}}(\xi_2, \mu_f) \right] \times \frac{d\hat{\sigma}(V_{\lambda_A} V'_{\lambda_B} \rightarrow \mathcal{F})}{dPS_n} \\ &+ \mathcal{O}\left(\frac{M_V^2}{M_{VV'}^2}\right) + \mathcal{O}\left(\frac{p_T^2}{M_{VV'}^2}\right) + \mathcal{O}\left(\frac{\alpha_W}{M_{VV'}^2} \log \frac{\mu_f^2}{M_{VV'}^2}\right). \end{aligned} \quad (2)$$

- Fully differential events for $V_{\lambda_A} V'_{\lambda_B} \rightarrow \mathcal{F}$ at lowest order (LO)
- $V_{\lambda_A} \in \{W_\lambda^\pm, Z_\lambda, \gamma_\lambda\}$ with nonzero M_Z, M_W
- \mathcal{F} can be anything (in practice, up to 5-6 legs)
- $f_{V_\lambda/\mu^\pm}(\xi, \mu_f)$ are bare LO PDFs [Dawson('85); Peskin&Schroeder]
- No RG evolution ← important!

W_λ/Z_λ PDFs depend on helicity (λ) of V and μ^\pm

- Subtle but important differences if evolving by q vs p_T

(this can account for some differences between groups!)

- γ_λ PDF (Weizsäcker-Williams) also added; kept “Improved WW”

$$f_{V+/f_L}(z, \mu_f^2) = \frac{g_V^2}{4\pi^2} \frac{g_L^2(1-z)^2}{2z} \log \left[\frac{\mu_f^2}{M_V^2} \right],$$

$$f_{V-/f_L}(z, \mu_f^2) = \frac{g_V^2}{4\pi^2} \frac{g_L^2}{2z} \log \left[\frac{\mu_f^2}{M_V^2} \right],$$

$$f_{V_0/f_L}(z, \mu_f^2) = \frac{g_V^2}{4\pi^2} \frac{g_L^2(1-z)}{z},$$

$$f_{V+/f_R}(z, \mu_f^2) = \left(\frac{g_R}{g_L} \right)^2 \times f_{V-/f_L}(z, \mu_f^2)$$

$$f_{V-/f_R}(z, \mu_f^2) = \left(\frac{g_R}{g_L} \right)^2 \times f_{V+/f_L}(z, \mu_f^2)$$

$$f_{V_0/f_R}(z, \mu_f^2) = \left(\frac{g_R}{g_L} \right)^2 \times f_{V_0/f_L}(z, \mu_f^2)$$

```

59 c  **** EVA (1/6) for f_L > v_+
60 c  double precision function eva_fL_to_vp(gg2,gL2,mv2,x,mu2,ievo)
61 implicit none
62 integer ievo          ! evolution by q2 or pT2
63 double precision gg2,gL2,mv2,x,mu2
64 double precision coup2,split,xxlog,fourPiSq
65 data fourPiSq/39.47841760435743d0/ ! = 4pi**2
66
67 c  print*, 'gg2,gL2,mv2,x,mu2,ievo',gg2 !3,gL2,mv2,x,mu2,ievo
68 coup2 = gg2*gL2/fourPiSq
69 split = (1.d0-x)**2 / 2.d0 / x
70 if(ievo.eq.0) then
71   xxlog = dlog(mu2/mv2)
72 else
73   xxlog = dlog(mu2/mv2/(1.d0-x))
74 endif
75
76 eva_fL_to_vp = coup2*split*xxlog
77 return
78 end
79
80 c  **** EVA (2/6) for f_L > v_-
81 c  double precision function eva_fL_to_vm(gg2,gL2,mv2,x,mu2,ievo)
82 implicit none
83 integer ievo          ! evolution by q2 or pT2
84 double precision gg2,gL2,mv2,x,mu2
85 double precision coup2,split,xxlog,fourPiSq
86 data fourPiSq/39.47841760435743d0/ ! = 4pi**2
87
88 coup2 = gg2*gL2/fourPiSq
89 split = 1.d0 / 2.d0 / x
90 if(ievo.eq.0) then
91   xxlog = dlog(mu2/mv2)
92 else
93   xxlog = dlog(mu2/mv2/(1.d0-x))
94 endif
95
96 eva_fL_to_vm = coup2*split*xxlog
97 return
98 end
99

```

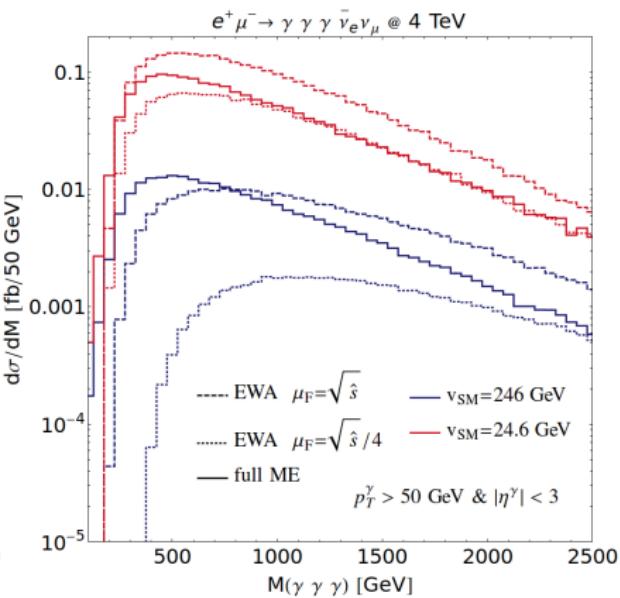
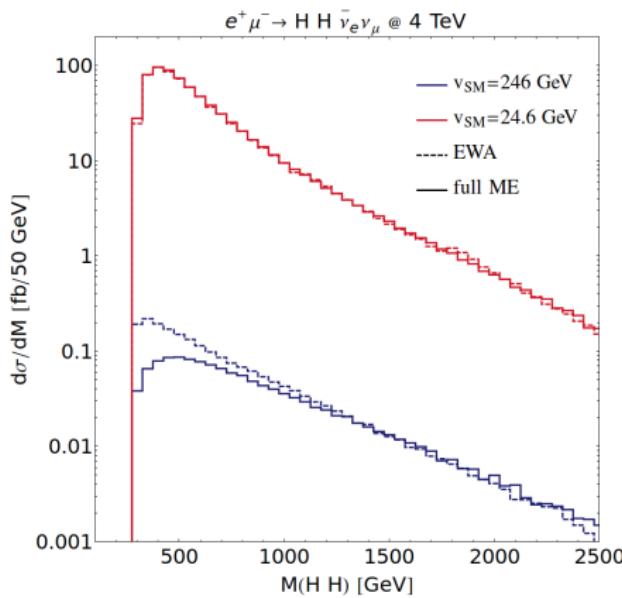
Universal and Quasi-Universal Uncertainties

Spent a lot of time trying to understand when factorization works

- Necessary that $(M_V^2/M_{VV'}^2) \lesssim 0.01$ (just like heavy quark factorization!)

(L) $W_0^+ W_0^- \rightarrow HH$

(R) $W_T^+ W_T^- \rightarrow \gamma\gamma\gamma$

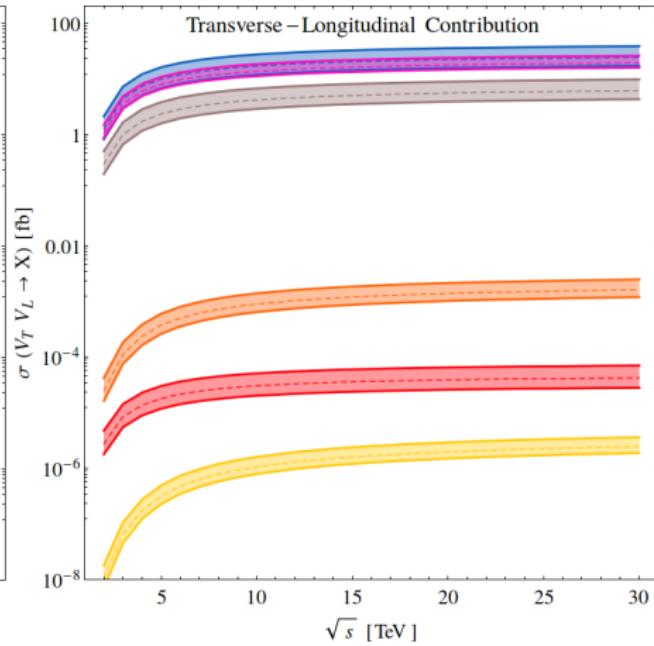
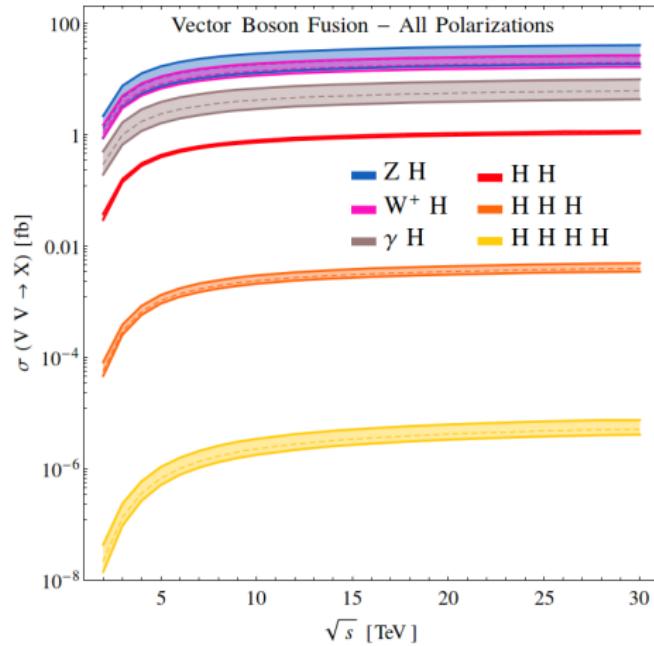


Survey (1/4)

We then had fun looking into *many* processes

$$(L) \sum_{\lambda_A, \lambda_B} V_{\lambda_A} V_{\lambda_B} \rightarrow HX$$

$$(R) V_T V_0 \rightarrow HX$$

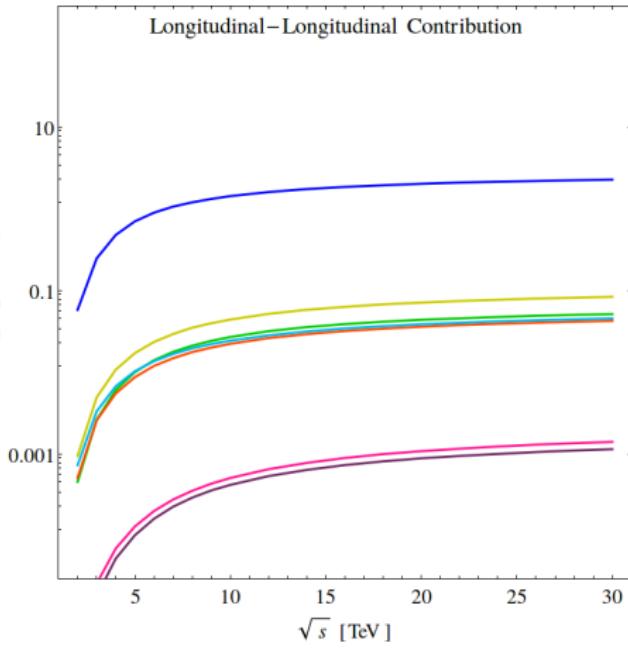
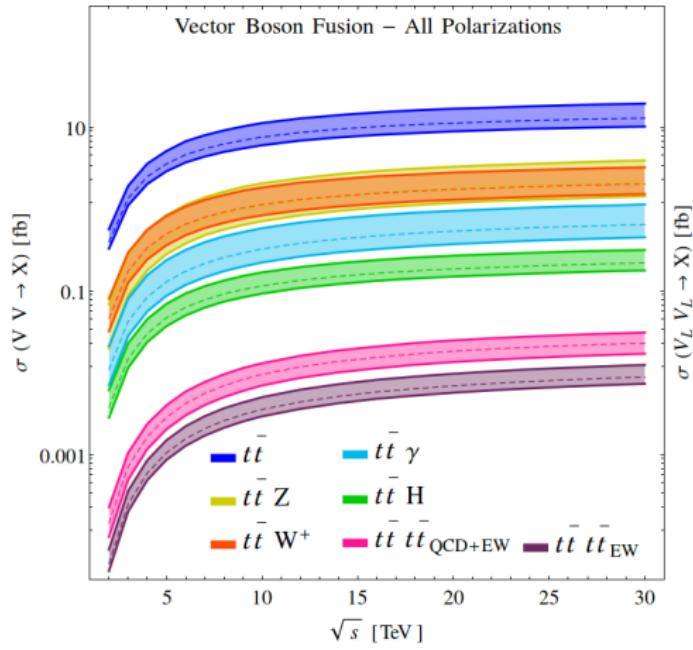


Survey (2/4)

We then had fun looking into *many* processes

$$(L) \sum_{\lambda_A, \lambda_B} V_{\lambda_A} V_{\lambda_B} \rightarrow t\bar{t}X$$

$$(R) V_0 V_0 \rightarrow t\bar{t}X$$

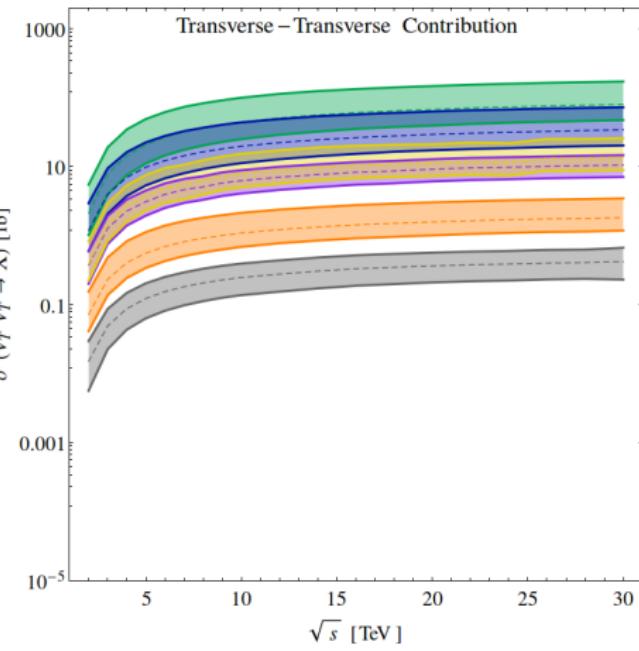
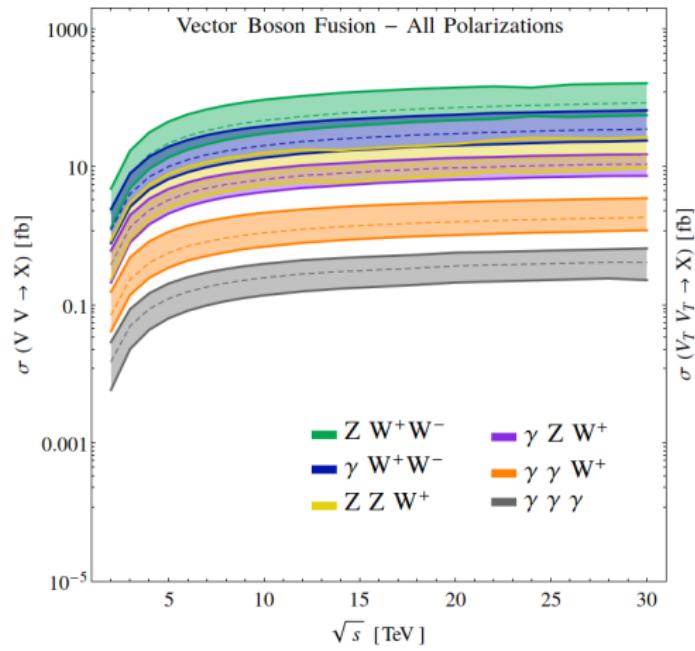


Survey (3/4)

We then had fun looking into *many* processes

$$(L) \sum_{\lambda_A, \lambda_B} V_{\lambda_A} V_{\lambda_B} \rightarrow 3V$$

$$(R) V_T V_T \rightarrow 3V$$



Survey (4/4)

We then had fun looking into *many* processes

	mg5amc syntax	$\sqrt{s} = 3 \text{ TeV}$	$\sqrt{s} = 14 \text{ TeV}$	$\sqrt{s} = 30 \text{ TeV}$
$\sum V_{\lambda_A} V'_{\lambda_B} \rightarrow W^+ W^-$	vp vm > w+ w-	$2.2 \cdot 10^2$ +98% -35%	$7.0 \cdot 10^2$ +91% -33%	$8.6 \cdot 10^2$ +88% -32%
$V_T V'_T \rightarrow W^+ W^-$	vp{T} vm{T} > w+ w-	$2.0 \cdot 10^2$ +99% -35%	$6.6 \cdot 10^2$ +93% -34%	$8.0 \cdot 10^2$ +92% -33%
$V_0 V'_T \rightarrow W^+ W^-$	vp{0} vm{T} > w+ w-	$1.2 \cdot 10^1$ +54% -27%	$4.4 \cdot 10^1$ +50% -25%	$5.2 \cdot 10^1$ +49% -24%
$V_0 V'_0 \rightarrow W^+ W^-$	vp{0} vm{0} > w+ w-	$4.2 \cdot 10^{-1}$	$1.7 \cdot 10^0$	$2.0 \cdot 10^0$
$\sum V_{\lambda_A} V'_{\lambda_B} \rightarrow W^+ Z$	vp vm > w+ z	$5.3 \cdot 10^1$ +105% -40%	$1.8 \cdot 10^2$ +97% -37%	$2.2 \cdot 10^2$ +95% -37%
$V_T V'_T \rightarrow W^+ Z$	vp{T} vm{T} > w+ z	$5.0 \cdot 10^1$ +111% -42%	$1.6 \cdot 10^2$ +103% -39%	$2.0 \cdot 10^2$ +100% -38%
$V_0 V'_T \rightarrow W^+ Z$	vp{0} vm{T} > w+ z	$3.4 \cdot 10^0$ +36% -18%	$1.4 \cdot 10^1$ +34% -17%	$1.7 \cdot 10^1$ +34% -17%
$V_0 V'_0 \rightarrow W^+ Z$	vp{0} vm{0} > w+ z	$3.9 \cdot 10^{-2}$	$2.1 \cdot 10^{-1}$	$2.6 \cdot 10^{-1}$
$\sum V_{\lambda_A} V'_{\lambda_B} \rightarrow ZZ$	vp vm > z z	$4.4 \cdot 10^1$ +164% -52%	$1.6 \cdot 10^2$ +144% -48%	$1.9 \cdot 10^2$ +143% -48%
$V_T V'_T \rightarrow ZZ$	vp{T} vm{T} > z z	$4.0 \cdot 10^1$ +171% -54%	$1.4 \cdot 10^2$ +153% -50%	$1.7 \cdot 10^2$ +150% -49%
$V_0 V'_T \rightarrow ZZ$	vp{0} vm{T} > z z	$4.2 \cdot 10^0$ +66% -33%	$1.8 \cdot 10^1$ +61% -30%	$2.2 \cdot 10^1$ +60% -30%
$V_0 V'_0 \rightarrow ZZ$	vp{0} vm{0} > z z	$1.1 \cdot 10^{-1}$	$6.0 \cdot 10^{-1}$	$7.2 \cdot 10^{-1}$
$\sum V_{\lambda_A} V'_{\lambda_B} \rightarrow \gamma Z$	vp vm > a z	$1.9 \cdot 10^1$ +169% -53%	$7.1 \cdot 10^1$ +149% -49%	$8.8 \cdot 10^1$ +145% -48%
$V_T V'_T \rightarrow \gamma Z$	vp{T} vm{T} > a z	$1.8 \cdot 10^1$ +172% -54%	$6.8 \cdot 10^1$ +153% -50%	$8.4 \cdot 10^1$ +149% -49%
$V_0 V'_T \rightarrow \gamma Z$	vp{0} vm{T} > a z	$9.5 \cdot 10^{-1}$ +67% -33%	$4.4 \cdot 10^0$ +61% -30%	$5.5 \cdot 10^0$ +60% -30%
$V_0 V'_0 \rightarrow \gamma Z$	vp{0} vm{0} > a z	$5.6 \cdot 10^{-4}$	$4.5 \cdot 10^{-3}$	$6.5 \cdot 10^{-3}$
$\sum V_{\lambda_A} V'_{\lambda_B} \rightarrow \gamma W^+$	vp vm > a w+	$1.1 \cdot 10^1$ +111% -42%	$4.0 \cdot 10^1$ +101% -39%	$4.9 \cdot 10^1$ +99% -38%
$V_T V'_T \rightarrow \gamma W^+$	vp{T} vm{T} > a w+	$1.1 \cdot 10^1$ +111% -42%	$3.9 \cdot 10^1$ +102% -39%	$4.8 \cdot 10^1$ +100% -38%
$V_0 V'_T \rightarrow \gamma W^+$	vp{0} vm{T} > a w+	$1.6 \cdot 10^{-2}$ +62% -31%	$7.3 \cdot 10^{-1}$ +56% -28%	$9.2 \cdot 10^{-1}$ +54% -27%
$V_0 V'_0 \rightarrow \gamma W^+$	vp{0} vm{0} > a w+	$1.5 \cdot 10^{-4}$	$1.2 \cdot 10^{-3}$	$1.7 \cdot 10^{-3}$
$\sum V_{\lambda_A} V'_{\lambda_B} \rightarrow \gamma\gamma$	vp vm > a a	$2.1 \cdot 10^0$ +172% -54%	$8.5 \cdot 10^0$ +152% -50%	$1.1 \cdot 10^1$ +147% -48%
$V_T V'_T \rightarrow \gamma\gamma$	vp{T} vm{T} > a a	$2.1 \cdot 10^0$ +172% -54%	$8.5 \cdot 10^0$ +152% -50%	$1.1 \cdot 10^1$ +147% -48%
$V_0 V'_T \rightarrow \gamma\gamma$	vp{0} vm{T} > a a	$7.8 \cdot 10^{-4}$ +70% -35%	$3.4 \cdot 10^{-3}$ +67% -34%	$4.2 \cdot 10^{-3}$ +67% -33%
$V_0 V'_0 \rightarrow \gamma\gamma$	vp{0} vm{0} > a a	$5.8 \cdot 10^{-4}$	$4.7 \cdot 10^{-3}$	$6.8 \cdot 10^{-3}$

How the EW theory behaves at very high energies is being explored!

- Exploring polarized matrix elements with MadGraph5

w/ D. Buarque Franzosi, O. Mattelaer, S. Shil [1912.01725]

- Exploring EW VBF with multi-TeV muon colliders

w/ A. Costantini, F. De Lillo, F. Maltoni, L. Mantani, O. Mattelaer, X. Zhao [2005.10289]

- ▶ Multi-TeV $\ell^+\ell^-$ colliders are “high-luminosity EW boson colliders”

- NEW! Exploring the validity of W_λ/Z_λ PDFs

w/ A. Costantini, F. Maltoni, O. Mattelaer [2111.02442]

- ▶ Support for W_λ/Z_λ PDFs from μ^\pm, e^\pm beams in MadGraph5



Thank you.

Example Usage

For $V_T V_T \rightarrow HH$ at $\sqrt{s} = 3$ TeV in $\mu^+ \mu^-$ collisions:

```
$ ./bin/mg5_aMC
set group_subprocesses false
set gauge Feynman
define vp = w+ z a
define vm = w- z a
generate vp{T} vm{T} > h h
output vtvt_hh
launch
set lpp1 -4
set lpp2 4
set pdlabel eva # calls EVA PDF sets
set fixed_fac_scale false
set dynamical_scale_choice 4 # muf = scalefact * dsqrt(shat)
set scalefact 0.5
set ievo_eva 0 # evolution scheme (0=q2 or 1=pT2)
set ebeam 1500 # GeV
set no_parton_cut
set nevents 1k
set dsqrt_shat 1000 # requirement on M(VV) system
set pt_min_pdg {25:50}
set eta_max_pdg {25:3}
set use_syst true
set nhel 1 # MC sampling over helicities. must be set to 1
done
```

See appendix B of [2111.02442] for many more details!